



Growth of cloud droplets from aerosol In turbulence

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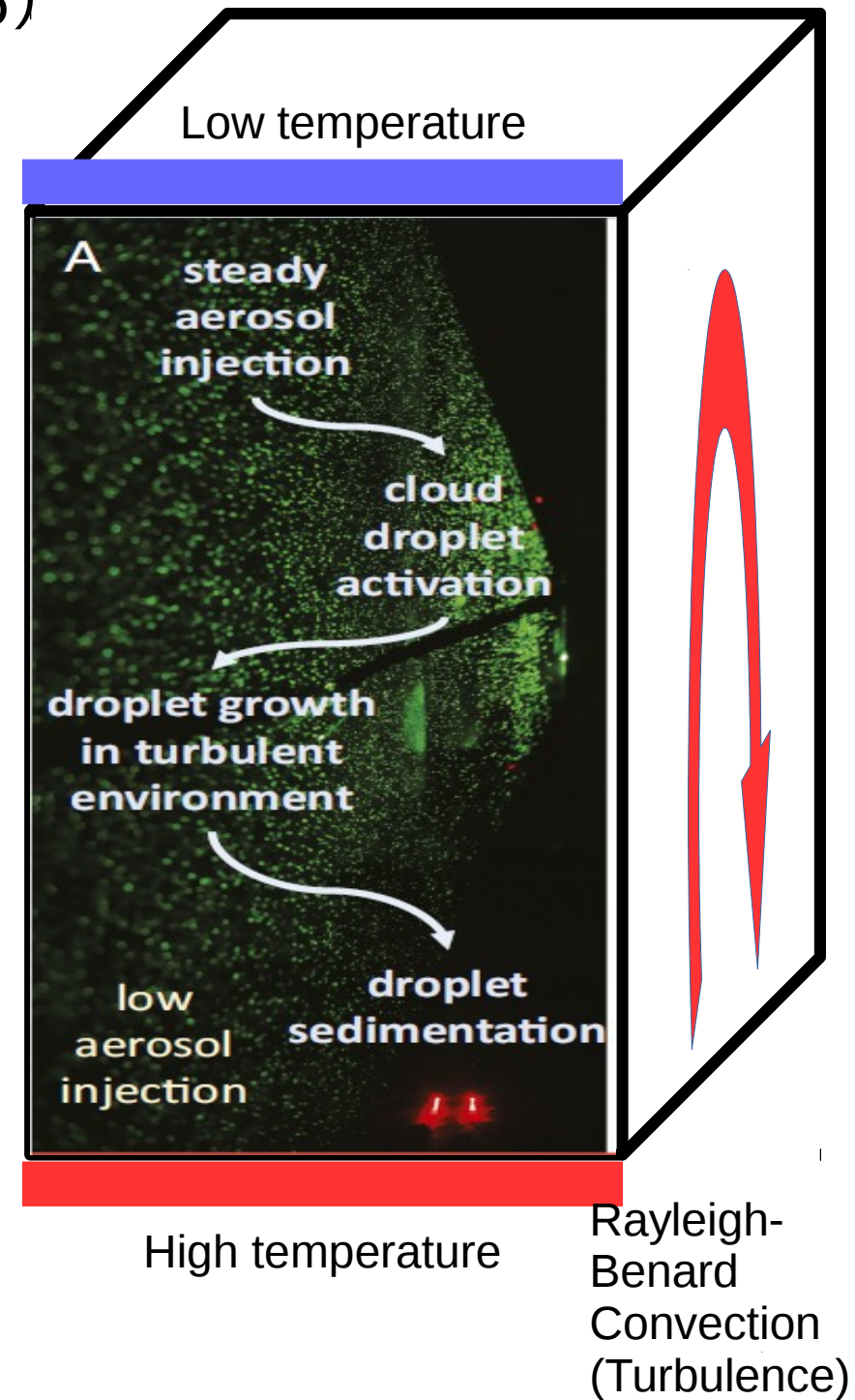
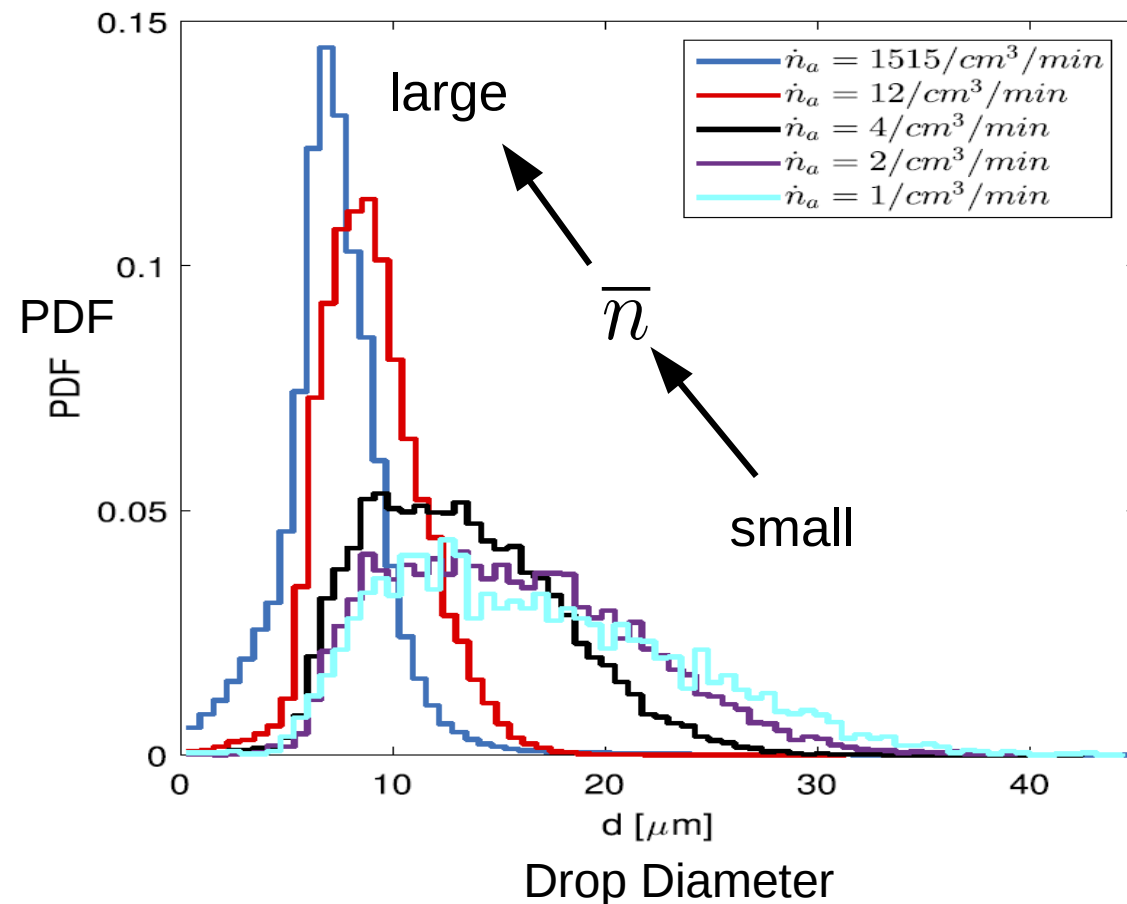
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Chandrakar et al (2016, PNAS) (C16)

- Moist R-B convection (turb.)
- Aerosol injection at const. rate
- Condensation nucleation
 - Condensation growth
 - Removal ($\Rightarrow$ steady state)



C16: Statistical theory and experimental results

Langevin Model

$$\frac{dr(t)^2}{dt} = 2KS$$

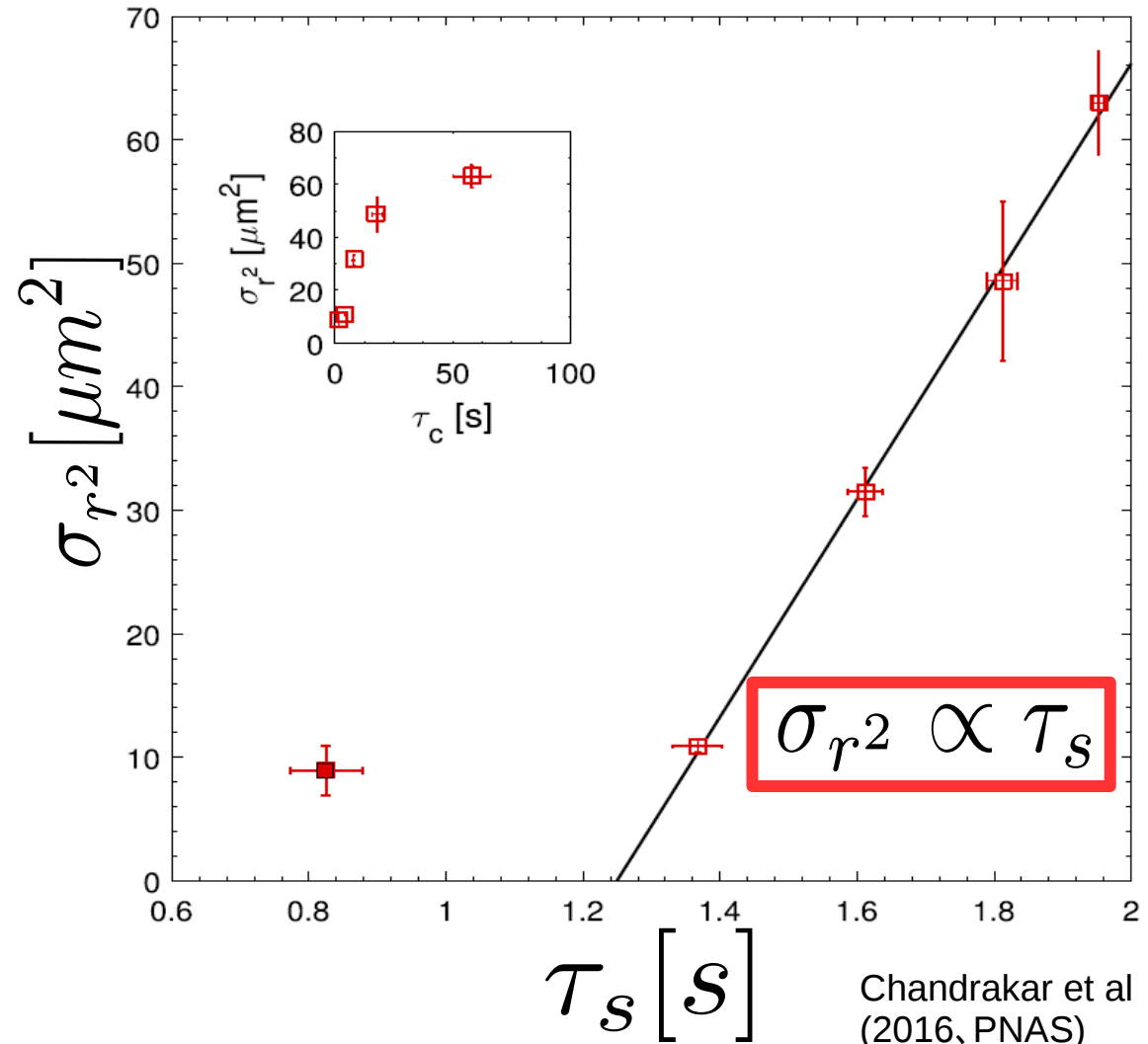
$$dS(t) = S(t + dt) - S(t)$$

$$= \left[\frac{S_{eq} - S}{\tau_t} - \frac{S}{\tau_c} \right] dt + \left(\frac{2\sigma_{s_0}^2 dt}{\tau_t} \right)^{1/2} \eta(t)$$



$$\overline{S'^2} = \left(\frac{\tau_s}{\tau_t} \right) \sigma_{s_0}^2$$

$$\sigma_{r^2} = \frac{8K^2 \sigma_{s_0}^2 \tau_s^2}{\tau_t} \tau_{res}$$



$$\overline{S'^2} \propto \tau_s$$

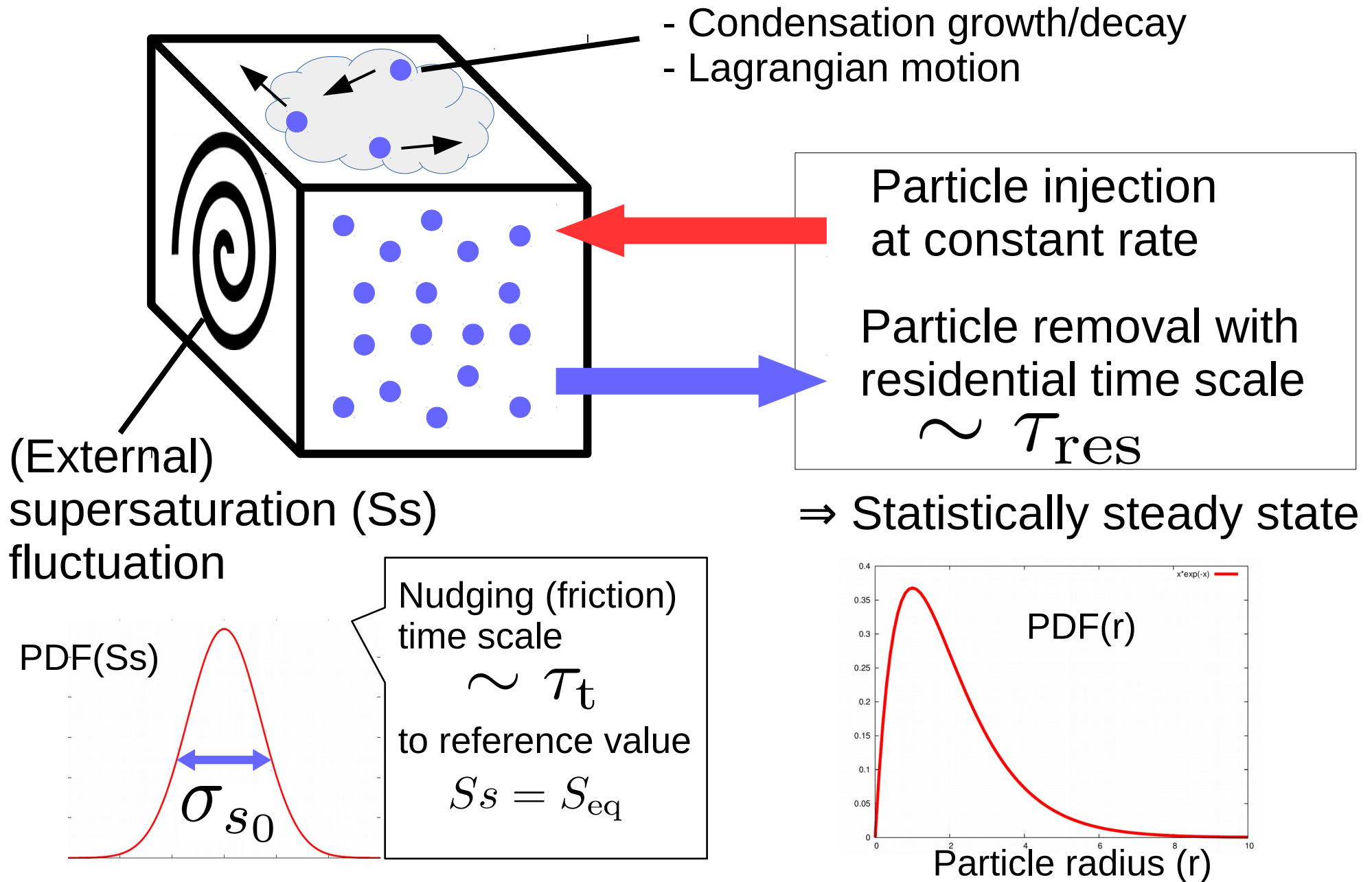
$$\sigma_{r^2} \propto \tau_s$$

System time scale

$$\left(\tau_s = \frac{\tau_c \tau_t}{\tau_c + \tau_t} \right)$$

phase change turb.

C16 (Langevin) model



Governing equations for DNS

Fluid

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\frac{1}{\rho_a} \nabla p + \nu_a \nabla^2 \mathbf{u} + \mathbf{f}_1$$

$$\frac{\partial T}{\partial t} + \mathbf{u} \cdot \nabla T = \frac{L_v}{c_p} C_d + \kappa_T \nabla^2 T - \frac{(T - T_{\text{eq}})}{\tau_t}$$

$$\frac{\partial Q}{\partial t} + \mathbf{u} \cdot \nabla Q = -C_d + \kappa_v \nabla^2 Q - \frac{(Q - Q_{\text{eq}})}{\tau_t} + \mathbf{f}_2$$

Nudging to
 $S_{\text{eq}}(T_{\text{eq}}, Q_{\text{eq}})$

Random force
($\Rightarrow$ Ss fluctuation)

Particle

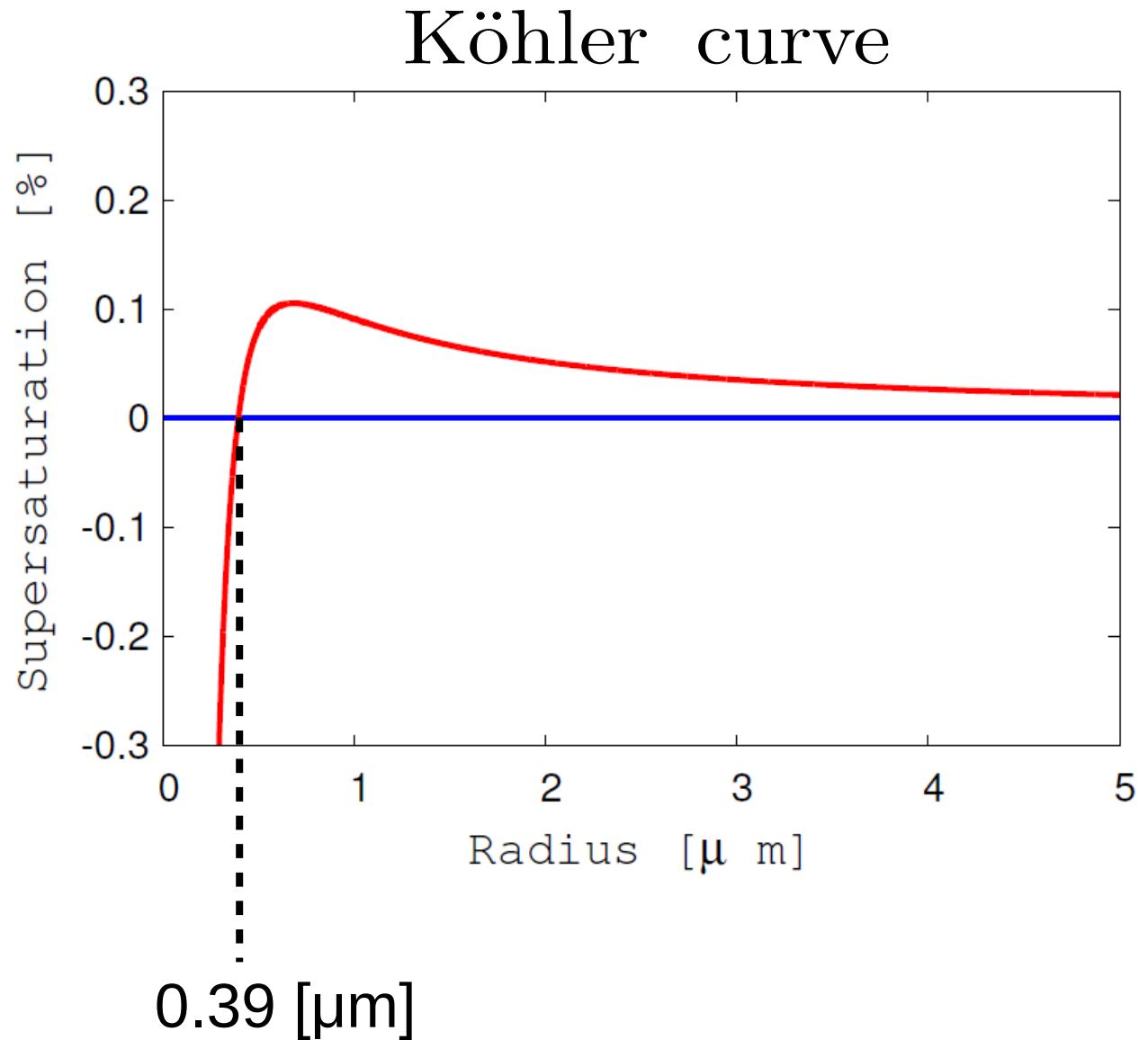
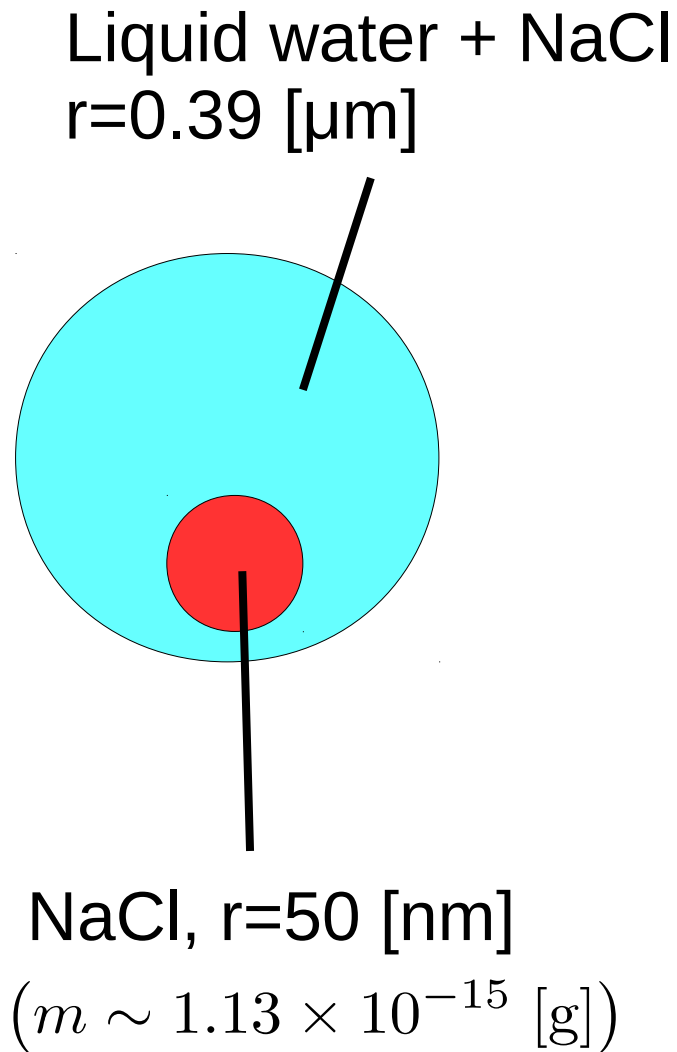
$$\frac{d\mathbf{X}_j(t)}{dt} = \mathbf{V}_j(t)$$

$$\frac{d\mathbf{V}_j(t)}{dt} = \frac{1}{\tau_j(t)} (\mathbf{u}(\mathbf{X}_j(t), t) - \mathbf{V}_j(t)) - \cancel{g} \mathbf{e}_z$$

$$\frac{dR_j(t)^2}{dt} = 2K_s(t)S(\mathbf{X}_j(t), t) + (\text{Köhler})$$

Aerosol effect

Injected aerosol particle: NaCl aq



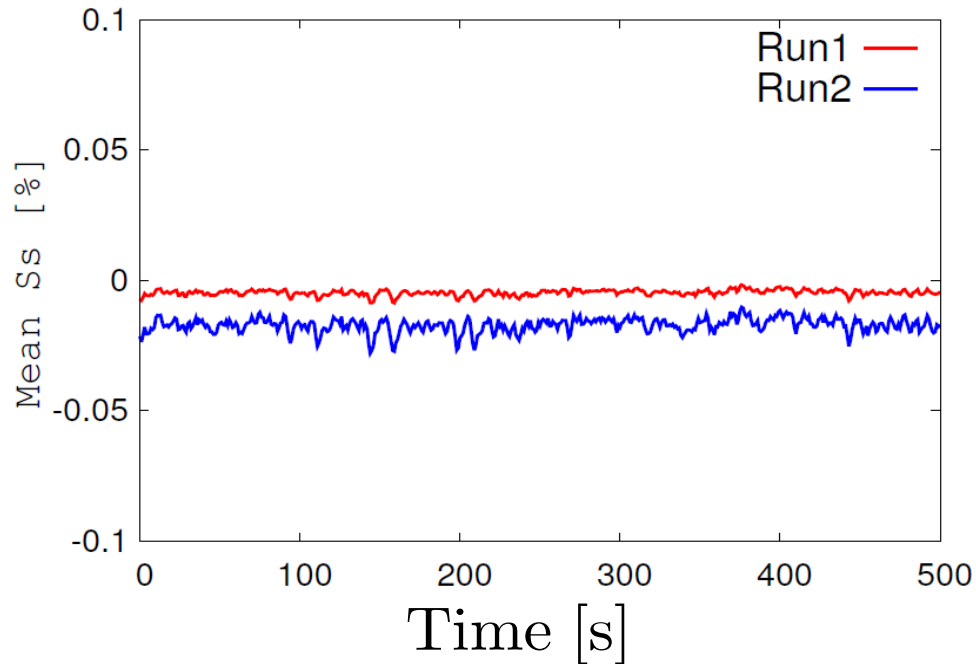
Experimental setups

Residential time scale	τ_{res}	606 [s] $\sim$ 10 [min]
Ss Nudging & fluctuation time scale	τ_{t}	2 [s]
Ss fluctuation without droplets	σ_{s_0}	1.4 [%] (Run1) 2.8 [%] (Run2)
Mean Ss in equilibrium	S_{eq}	<u>0 [%]</u>

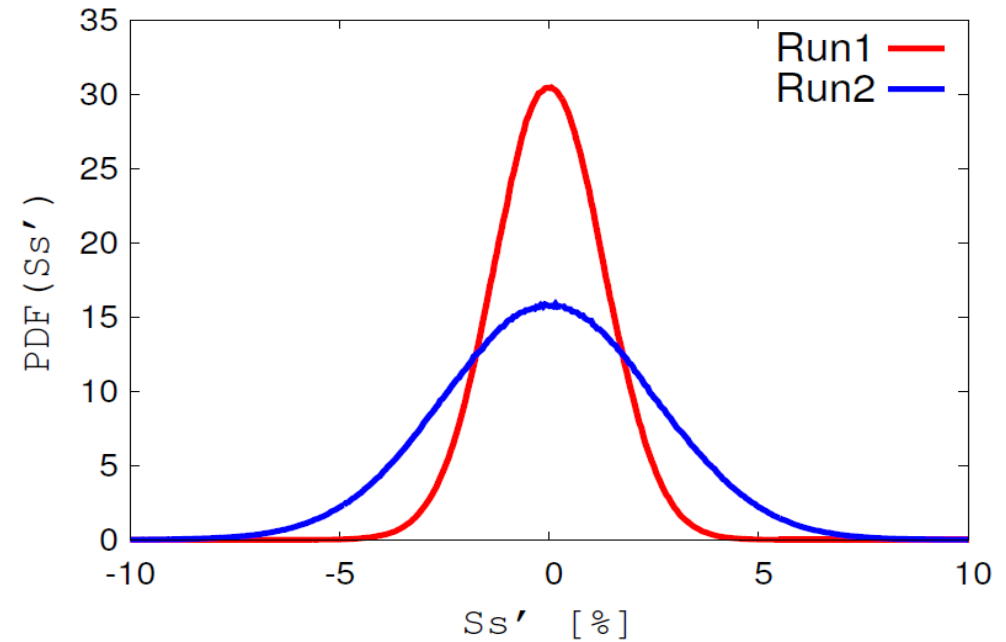
Results

Supersaturation

Volume ave. $Ss(t)$ [%]



PDF of Ss' [%]

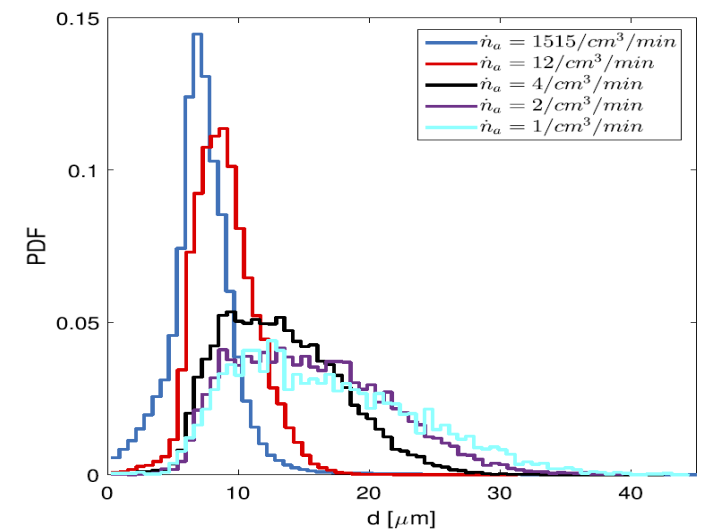
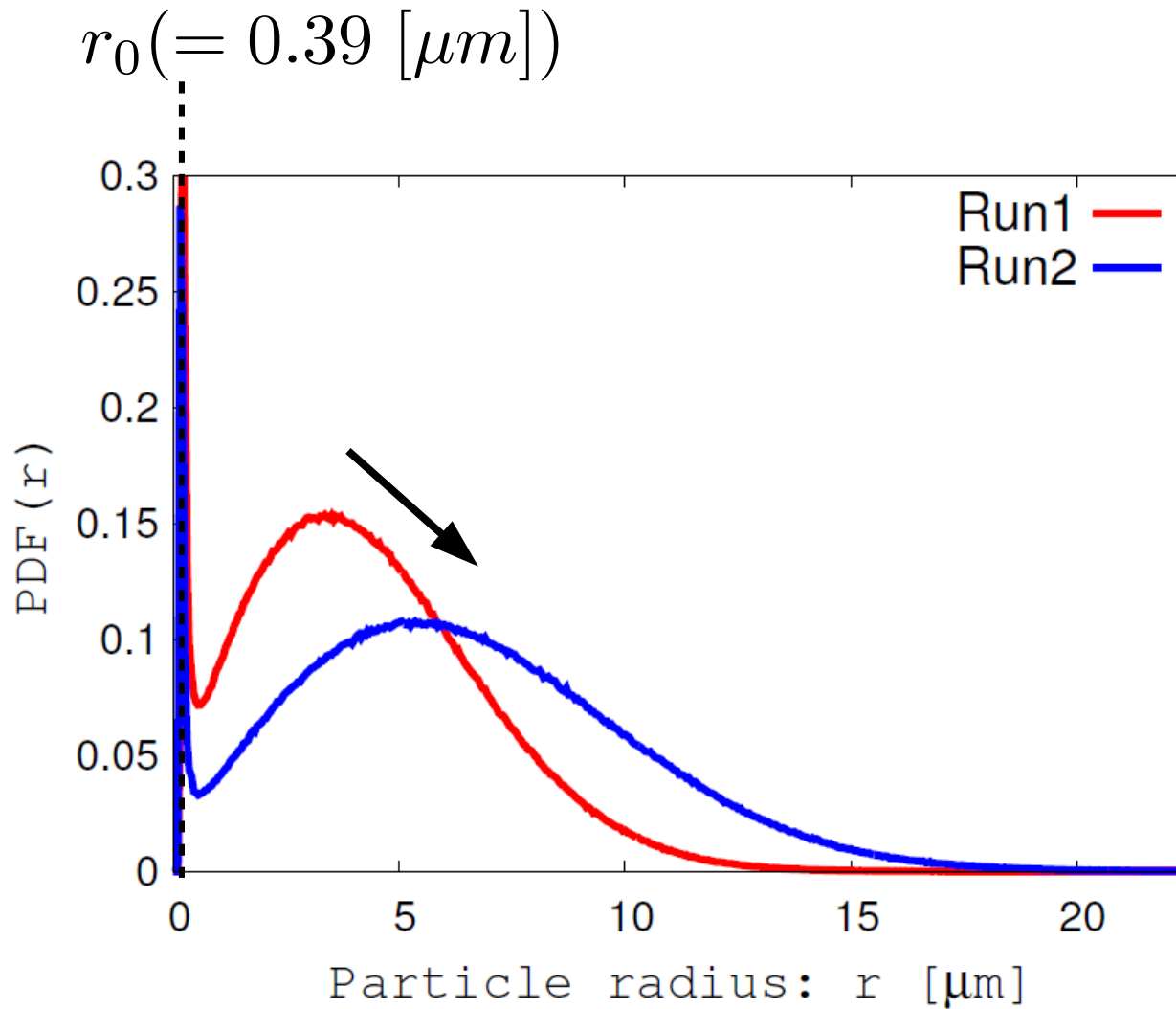


$$\bar{n} \sim 63 \text{ [cm}^{-3}\text{]}$$

$$\tau_c \sim 12.5 \text{ [s]} \quad \tau_t = 2 \text{ [s]}$$

$$\tau_s = \frac{\tau_t \tau_c}{\tau_t + \tau_c} \sim 1.7 \text{ [s]}$$

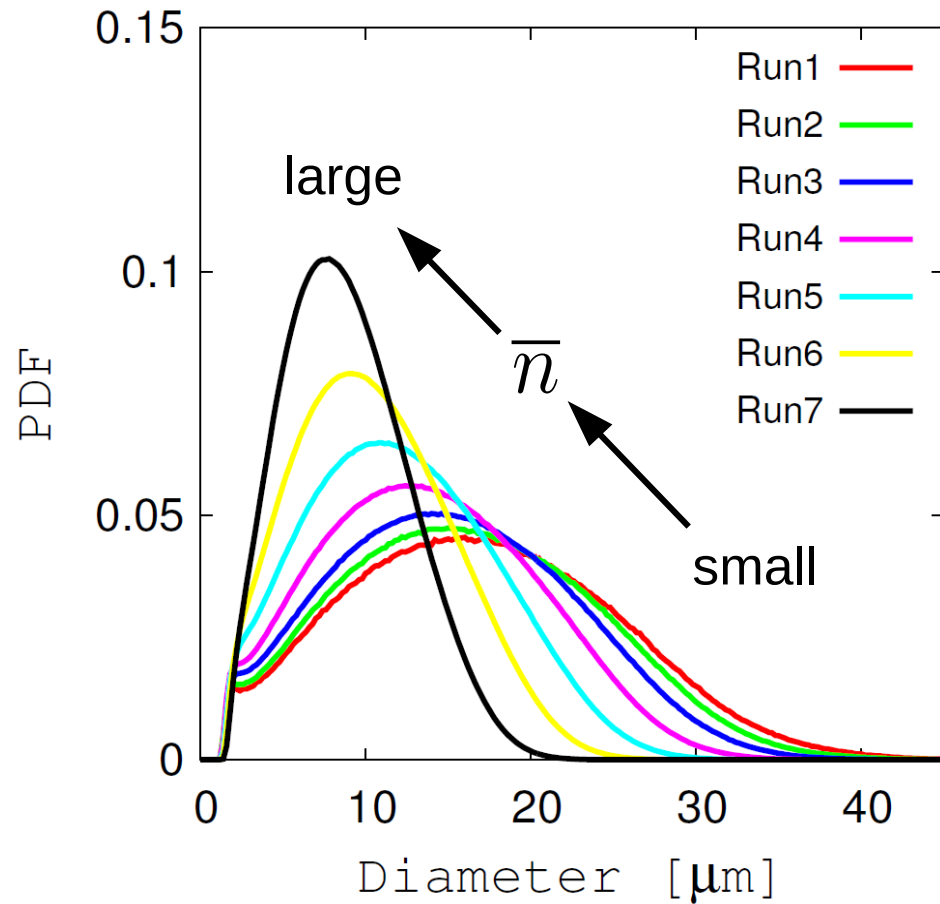
Size distribution (steady state)



(C16, Fig.2)

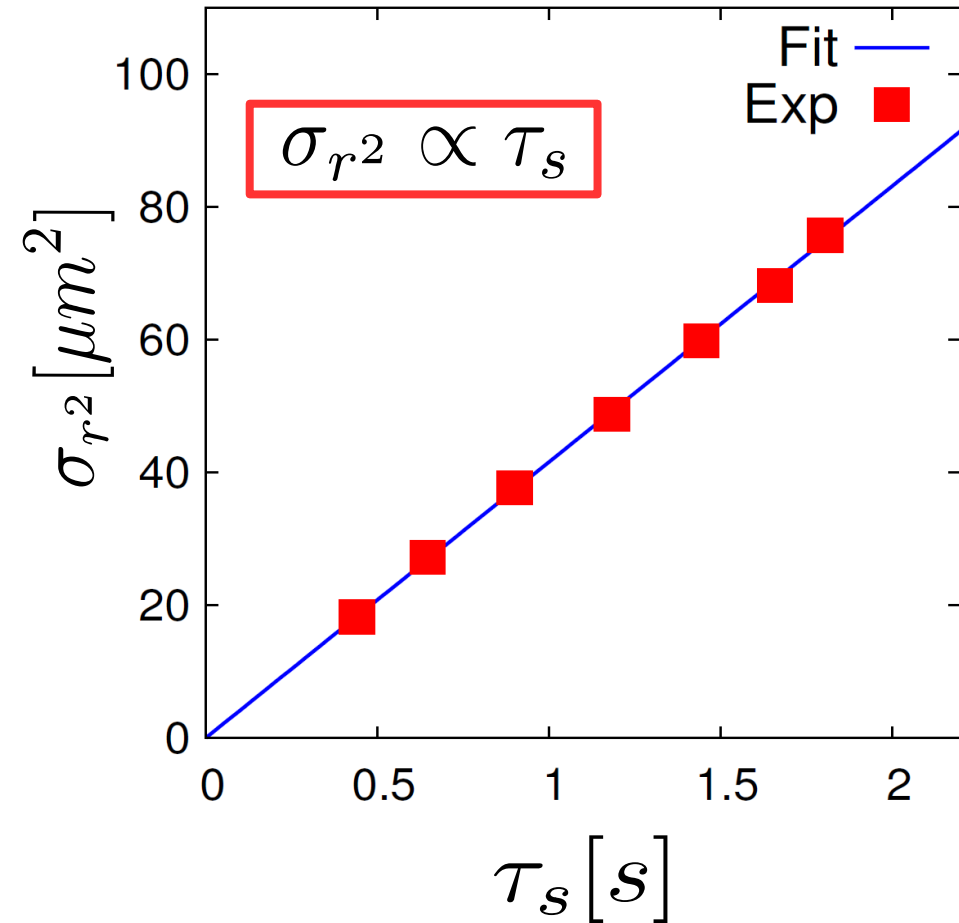
(Results from other experiments)

PDF of droplet diameter

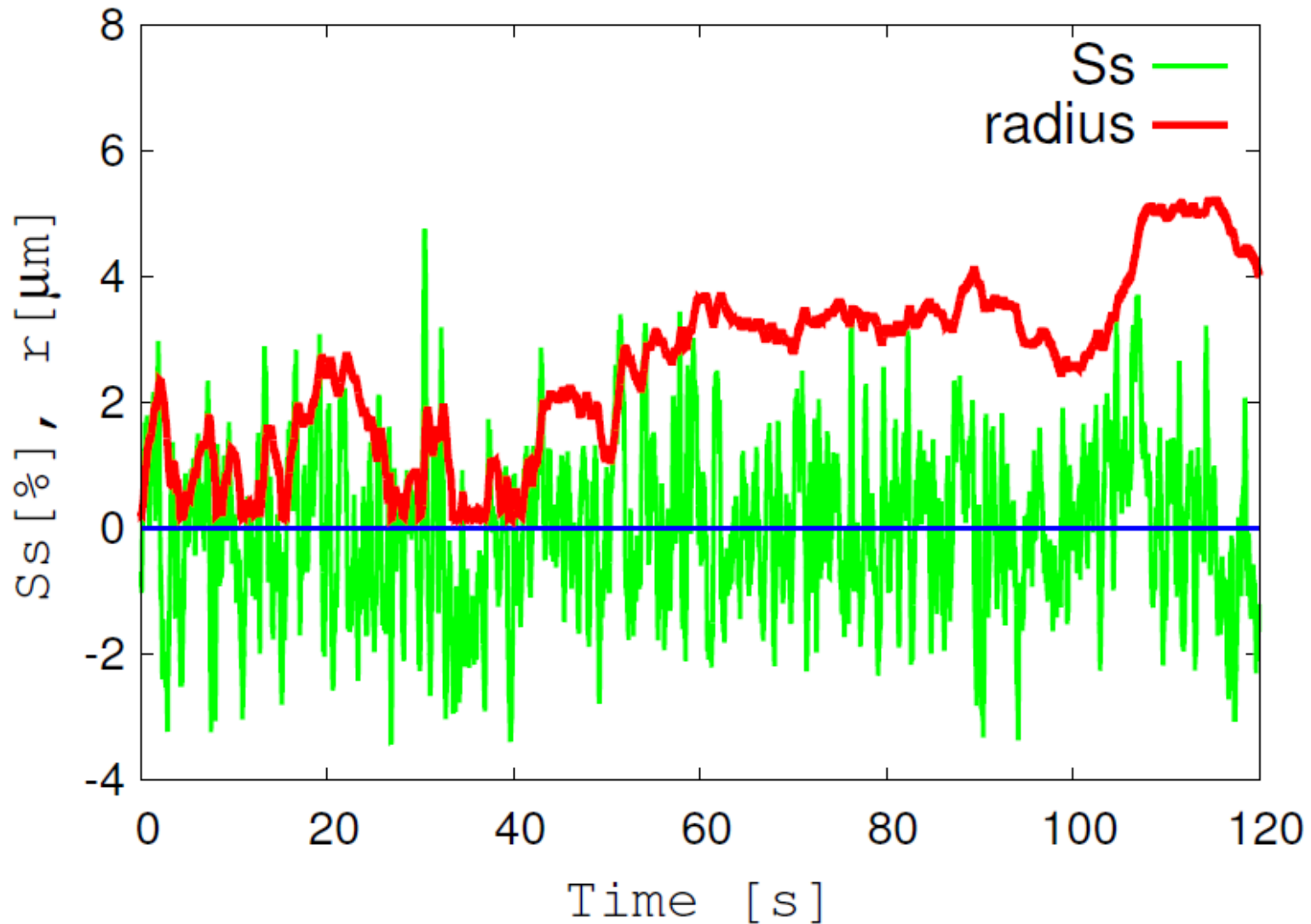


$$(S_{\text{eq}} = 0.1 [\%])$$

σ_{r^2} VS τ_s

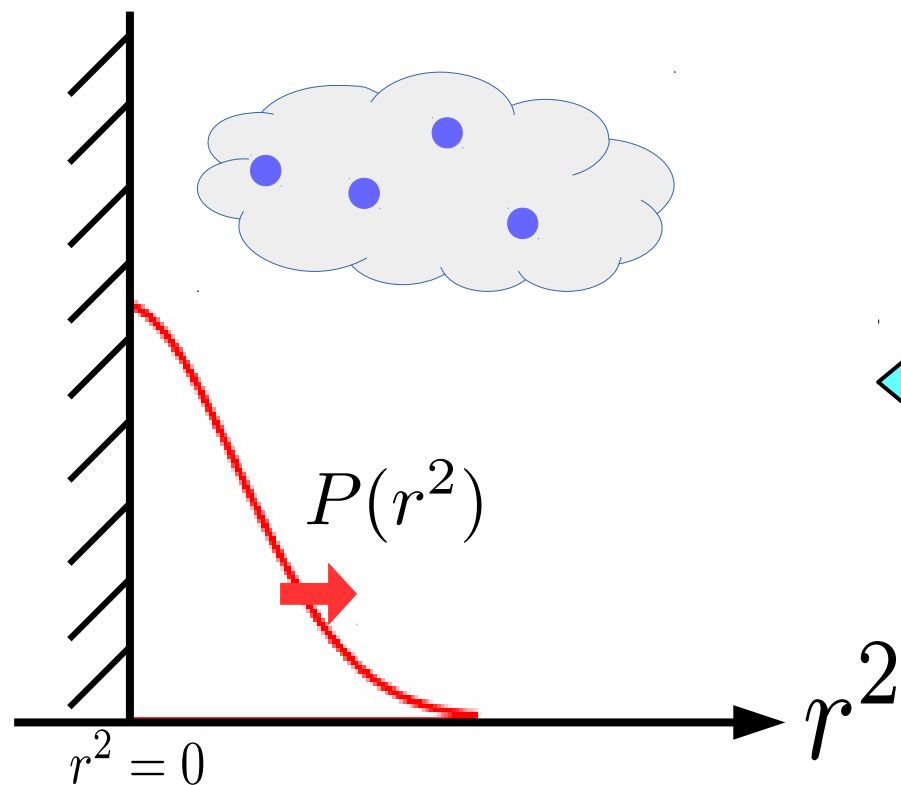


Run1: Lagrangian tracking, Ex1



←→ relatively $Ss < 0$ ("unlucky") ←→ relatively $Ss > 0$ ("lucky")

Droplet growth in fluctuating S s
with aerosol effect

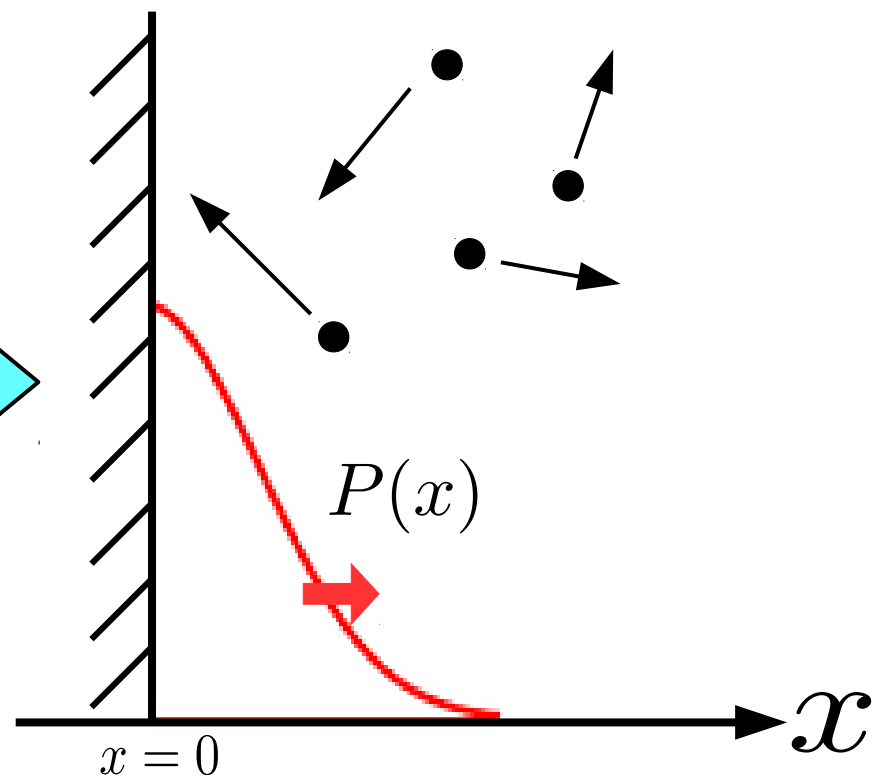


$$\frac{d\boxed{r^2}}{dt} = 2KS + (\text{Köhler})$$

$$\frac{d\boxed{S}}{dt} = -\frac{S}{\tau} + f$$

$(r^2 > 0)$

Brownian motion
with reflecting wall



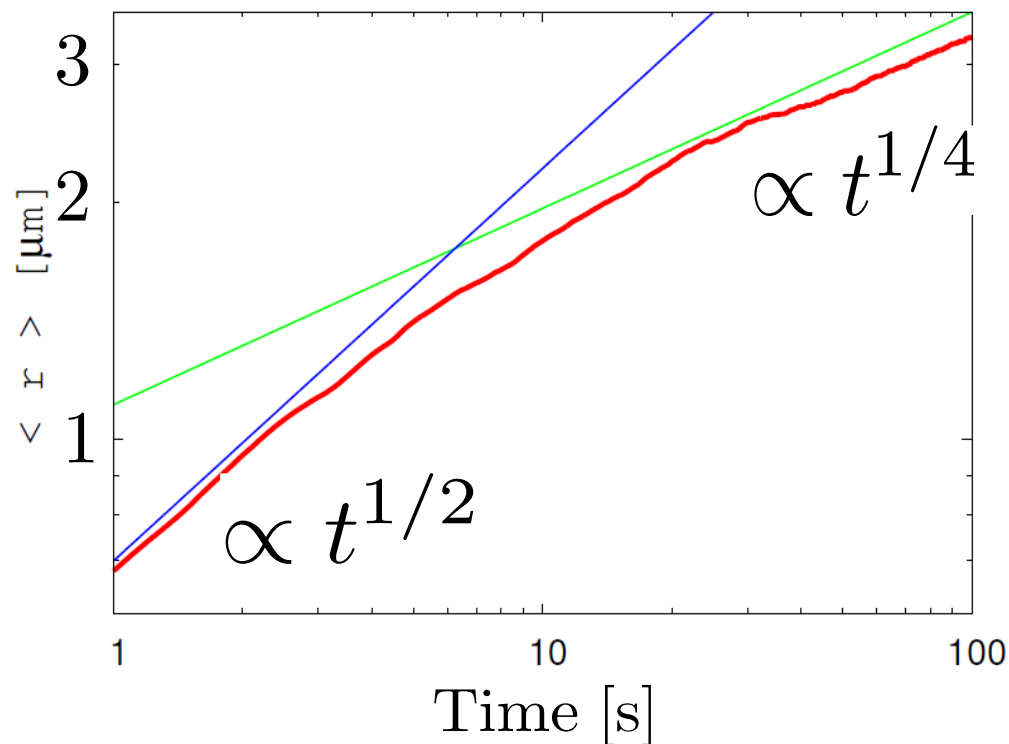
$$\frac{d\boxed{x}}{dt} = v$$

$$\frac{d\boxed{v}}{dt} = -\frac{v}{\tau} + f$$

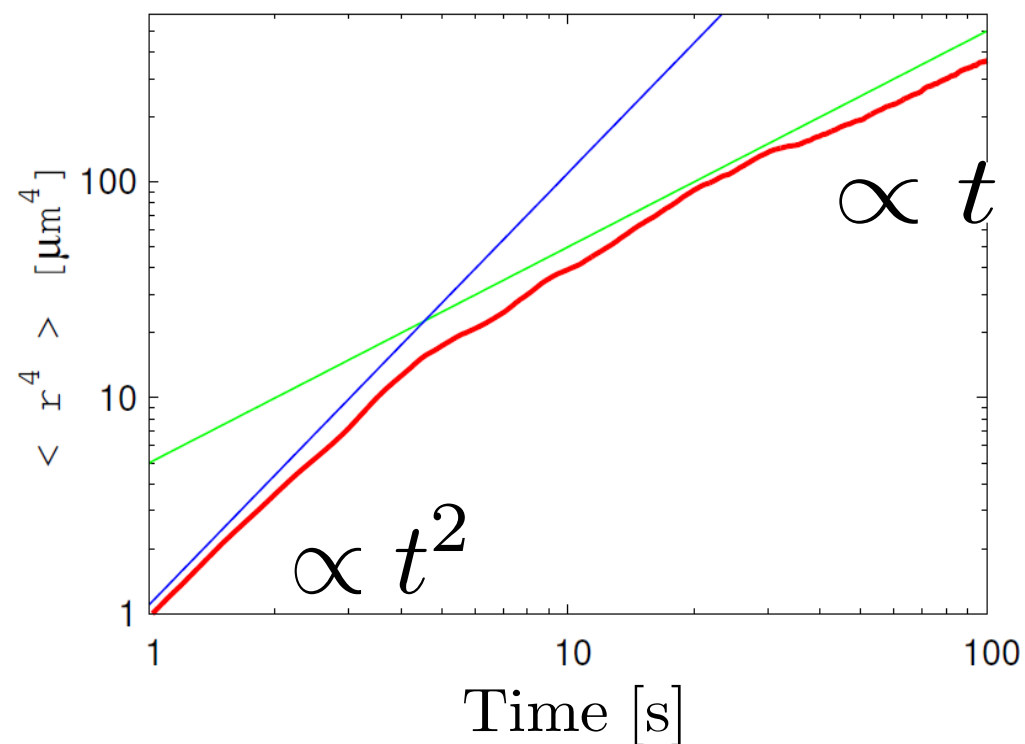
$(x > 0)$

Droplets	Brownian motion with wall	Scaling
$\langle r^4 \rangle$	$\langle x^2 \rangle$	$\sim Dt$ (C16)
$\langle r \rangle$	$\langle \sqrt{x} \rangle$	$\sim (Dt)^{1/4}$

$\langle r \rangle$ (Run1) (Red)



$\langle r^4 \rangle$ (Run1) (Red)



Summary

- Experimental results by Chandrakar et al (2016) was reproduced well by DNS
- Increase of mean radius due to aerosol effect (= wall effect, Siewert et al, 2017) in experiments with $Seq=0$ was confirmed
- This mechanism might play important role in C16's chamber experiments, (and possibly in real clouds? where mean S_s is small).

(typically $\overline{S_s} \sim 0.1$ [%] in clouds)

(Siebert & Shaw, 2017, JAS)